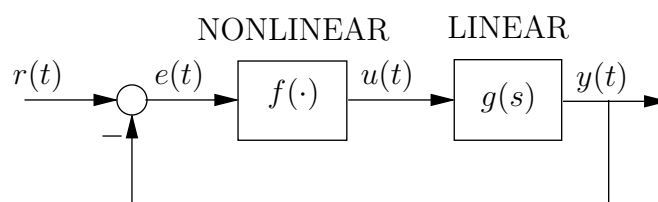


## Handout 3: Describing Functions

### 1 Harmonic balance

The *describing function method* (also called the method of harmonic balance) uses frequency domain (Fourier series) techniques to investigate limit cycle behaviour in nonlinear systems. The method involves an approximation, but nevertheless often gives a reliable prediction concerning limit cycle behaviour.

The usual context for the method is the following feedback system.



There is a linear system with transfer function  $g(s)$  and a memoryless nonlinearity  $f(\cdot)$  (e.g. saturation, hysteresis, backlash). For simplicity we will take the case of  $r = 0$ . We are looking for possible limit cycle (i.e. periodic) behaviour in the feedback system. Therefore let us take a trial solution:

$$e(t) = E \sin(\omega t).$$

If the output of the non-linearity is periodic with frequency  $\omega$  there will be a Fourier series representation:

$$u(t) = U_0 + \sum_{k=1}^{\infty} (U_k \sin(k\omega t) + V_k \cos(k\omega t)).$$

The Fourier coefficients corresponding to the fundamental harmonic are given by

$$U_1 = \frac{\omega}{\pi} \int_0^{2\pi/\omega} f(E \sin \omega t) \sin \omega t dt = \frac{1}{\pi} \int_0^{2\pi} f(E \sin \theta) \sin \theta d\theta,$$

$$V_1 = \frac{1}{\pi} \int_0^{2\pi} f(E \sin \theta) \cos \theta d\theta.$$

We now consider the approximation:

$$u(t) \approx U_1 \sin(\omega t) + V_1 \cos(\omega t)$$

as the input to the linear system. In the steady state, the output of the linear system is:

$$\begin{aligned} y(t) &= |g(j\omega)|[U_1 \sin(\omega t + \phi) + V_1 \cos(\omega t + \phi)] \\ &= \text{Im}[|g(j\omega)|(U_1 + jV_1)e^{j\omega t + \phi}] \\ &= \text{Im}[g(j\omega)(U_1 + jV_1)e^{j\omega t}] \end{aligned}$$

where  $\phi = \arg(g(j\omega))$ . (We remark that the neglect of the higher harmonics is most likely to be a valid approximation when  $g(s)$  behaves like a low pass filter.) In order to satisfy the feedback equations we need to have  $y(t) = -e(t)$ . We thus obtain

$$\begin{aligned} 0 &\equiv y(t) + e(t) \\ \Rightarrow 0 &\equiv \text{Im}[(g(j\omega)(U_1 + jV_1) + E)e^{j\omega t}] = 0 \\ \Rightarrow 0 &= g(j\omega)(U_1 + jV_1) + E. \end{aligned}$$

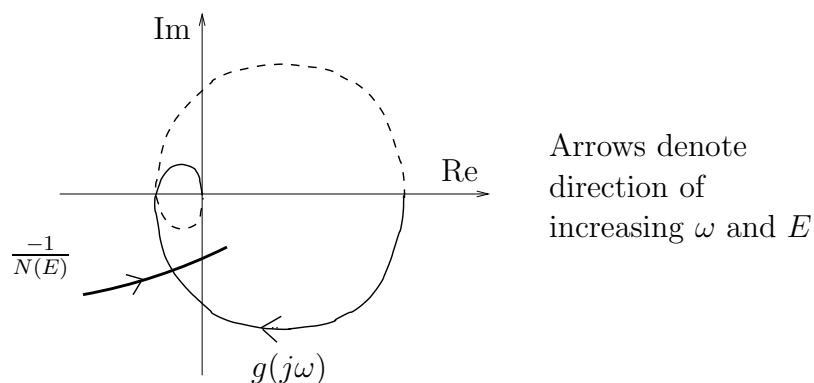
This is the condition for “harmonic balance”, i.e. for the fundamental harmonic to solve the feedback equations. We now define the *describing function*

$$N(E) = \frac{U_1 + jV_1}{E}.$$

Then the condition for harmonic balance becomes

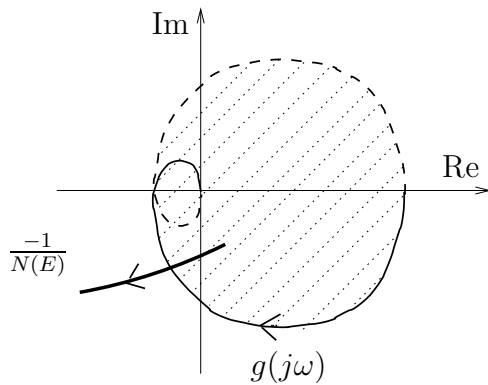
$$g(j\omega) = \frac{-1}{N(E)}.$$

This condition has an interesting graphical interpretation. Let the Nyquist diagram  $g(j\omega)$  of the linear system be plotted in the complex plane. On the same diagram plot the locus

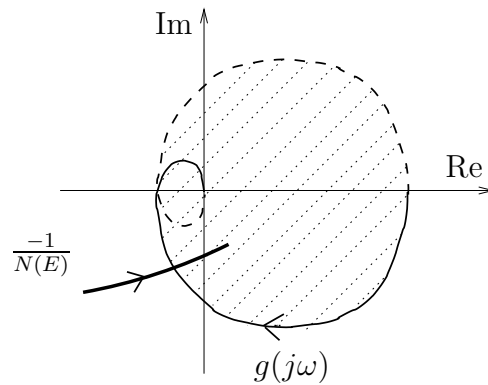


$-1/N(E)$ . Note that these loci can be plotted independently since they are respectively functions of  $\omega$  and  $E$  only. A limit cycle is predicted if there is an intersection of the two loci. Moreover, the intersection point gives an estimate for the frequency  $\omega$  and amplitude  $E$  of the limit cycle.

A useful prediction of the stability of limit cycles can be made from the manner of intersection of the two loci. Recall from the Nyquist stability criterion that, for an *open* loop stable  $g(s)$  in the standard unity gain negative feedback configuration (non-linearity absent of course), the *closed* loop feedback system is stable providing the critical point  $(-1, 0)$  lies to the left of the Nyquist locus. The point  $-1/N(E)$  can be thought of as the ‘critical point’ with respect to the Nyquist locus. Thus, if a small increase of the amplitude  $E$  moves the value of  $-1/N(E)$  *outside* of the shaded region below, then a *stabilizing* effect will be predicted, which will act to reduce the amplitude. Conversely, if an increase of  $E$  moves  $-1/N(E)$  *inside* the shaded region below, then a *destabilizing* effect will be predicted.



PREDICT STABLE LIMIT CYCLE



PREDICT UNSTABLE LIMIT CYCLE

## 2 Calculation of describing functions

### 2.1 Relay nonlinearity

$f(e) = \text{sign}(e)$ . Since  $\text{sign}(\sin(\omega t))$  is an odd function then  $V_1 = 0$ . Also

$$\begin{aligned} U_1 &= \frac{1}{\pi} \int_0^{2\pi} \text{sign}(E \sin \theta) \sin \theta \, d\theta, \\ &= \frac{1}{\pi} \int_0^{\pi} \sin \theta \, d\theta - \frac{1}{\pi} \int_{\pi}^{2\pi} \sin \theta \, d\theta \\ &= \frac{4}{\pi}. \end{aligned}$$

Thus  $N(E) = \frac{4}{\pi E}$ .

## 2.2 Polynomial nonlinearity

$f(e) = e^n$ . If  $n$  is odd it can be shown (from De Moivre's Theorem and the binomial expansion) that

$$\sin^n \theta = \frac{2}{2^n} \sum_{k=0}^{\frac{n-1}{2}} (-1)^{\binom{n-1}{2}-k} \binom{n}{k} \sin([n-2k]\theta)$$

This is a Fourier series with a finite number of terms; taking the last term,  $k = \frac{n-1}{2}$ , gives us the fundamental component if  $e = (E \sin \theta)^n$ :

$$U_1 \sin \theta = E^n \frac{2}{2^n} (-1)^0 \binom{n}{\frac{n-1}{2}} \sin \theta, \quad V_1 = 0.$$

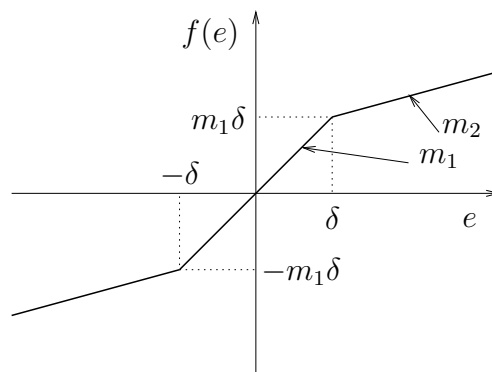
A familiar case of this is  $n = 3$  (eg in the Maths Databook):

$$(E \sin \theta)^3 = E^3 \left( \frac{3}{4} \sin \theta - \frac{1}{4} \sin(3\theta) \right)$$

for which  $U_1 = \frac{3E^3}{4}$ ,  $V_1 = 0$ , so  $N(E) = \frac{3E^2}{4}$ .

If  $n$  is even a similar formula exists, but the fundamental component is always zero (for example,  $\sin^2 \theta = \frac{1}{2} - \frac{1}{2} \cos(2\theta)$ ), so our low-pass assumption will eliminate everything. There is also a non-zero constant component — which can be handled by an extension of the harmonic balance idea.

## 2.3 Piecewise-linear non-linearity



$$f(e) = \begin{cases} m_1 e & \text{if } |e| < \delta \\ (m_1 - m_2)\delta + m_2 e & \text{if } e > \delta \\ (m_2 - m_1)\delta + m_2 e & \text{if } e < -\delta \end{cases}$$

Once again, since  $f(e)$  is an odd function  $V_1 = 0$ . If  $E \leq \delta$  then

$$U_1 = \frac{1}{2\pi} \int_0^\pi m_1 E \sin^2 \theta d\theta = m_1 E.$$

If  $E > \delta$  then

$$\begin{aligned} U_1 &= \frac{4}{\pi} \int_0^{\pi/2} f(E \sin \theta) \sin \theta d\theta, \\ &= \frac{4}{\pi} \int_0^{\sin^{-1}(\delta/E)} m_1 E \sin^2 \theta d\theta + \frac{4}{\pi} \int_{\sin^{-1}(\delta/E)}^{\pi/2} ((m_1 - m_2)\delta + m_2 E \sin \theta) \sin \theta d\theta \\ &= \frac{2m_1 E}{\pi} \left[ \theta - \frac{1}{2} \sin(2\theta) \right]_0^{\sin^{-1}(\delta/E)} + \frac{4(m_2 - m_1)\delta}{\pi} [\cos \theta]_{\sin^{-1}(\delta/E)}^{\pi/2} \\ &\quad + \frac{2m_2 E}{\pi} \left[ \theta - \frac{1}{2} \sin(2\theta) \right]_{\sin^{-1}(\delta/E)}^{\pi/2}. \end{aligned}$$

Now, since  $\cos(\sin^{-1} x) = \sqrt{1 - x^2}$  and  $\sin(2 \sin^{-1} x) = 2x\sqrt{1 - x^2}$  for  $0 < x < \pi/2$ , we get after a little algebra:

$$U_1 = \frac{2(m_1 - m_2)E}{\pi} \left[ \sin^{-1}\left(\frac{\delta}{E}\right) + \frac{\delta}{E} \left(1 - \left(\frac{\delta}{E}\right)^2\right)^{1/2} \right] + m_2 E.$$

Hence we get:

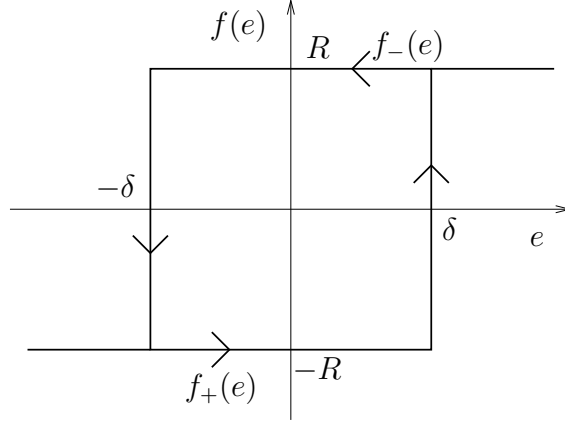
$$N(E) = \begin{cases} m_1, & \text{if } E < \delta \\ \frac{2(m_1 - m_2)}{\pi} \left[ \sin^{-1}\left(\frac{\delta}{E}\right) + \frac{\delta}{E} \left(1 - \left(\frac{\delta}{E}\right)^2\right)^{1/2} \right] + m_2, & \text{if } E > \delta. \end{cases}$$

## 2.4 Relay with hysteresis

Consider the following two-valued function:

This is a non-linearity with ‘memory’.  $f(e)$  takes the value  $+R$  or  $-R$  according to whether  $e$  was greater than  $\delta$  or less than  $-\delta$  on the last occasion when  $|e| > \delta$ . We will find the describing function for  $E > \delta$  since it is not properly defined otherwise. We have

$$N(E) = \frac{1}{\pi E} \int_{-\pi/2}^{3\pi/2} f(E \sin \theta) (\sin \theta + j \cos \theta) d\theta$$



since we can integrate over any cycle of length  $2\pi$ . Thus

$$N(E) = \frac{1}{\pi E} \int_{-\pi/2}^{\pi/2} f_+(E \sin \theta) (\sin \theta + j \cos \theta) d\theta + \frac{1}{\pi E} \int_{\pi/2}^{3\pi/2} f_-(E \sin \theta) (\sin \theta + j \cos \theta) d\theta.$$

Let  $\nu = E \sin \theta$ ,  $d\nu = E \cos \theta d\theta$  (and note that  $\cos \theta$  has positive sign in the first integral and negative sign in the second), then

$$\begin{aligned} \operatorname{Re}N(E) &= \frac{1}{\pi E} \int_{-\pi/2}^{\pi/2} f_+(E \sin \theta) \sin \theta d\theta + \frac{1}{\pi E} \int_{\pi/2}^{3\pi/2} f_-(E \sin \theta) \sin \theta d\theta \\ &= \frac{1}{\pi E} \int_{-E}^E f_+(\nu) \left(\frac{\nu}{E}\right) \frac{d\nu}{E\sqrt{1-(\nu/E)^2}} - \frac{1}{\pi E} \int_E^{-E} f_-(\nu) \left(\frac{\nu}{E}\right) \frac{d\nu}{E\sqrt{1-(\nu/E)^2}} \\ &= \frac{1}{\pi E} \int_{-E}^E [f_+(\nu) + f_-(\nu)] \left(\frac{\nu}{E}\right) \frac{d\nu}{E\sqrt{1-(\nu/E)^2}}. \end{aligned}$$

The above formula is valid for any two-valued non-linearity. In fact, it is also valid for a single-valued non-linearity  $f$ . If  $f = (f_+ + f_-)/2$  is an odd function then

$$N(E) = \frac{1}{\pi E} \int_{-E}^E 2f(\nu) \left(\frac{\nu}{E}\right) \frac{d\nu}{E\sqrt{1-(\nu/E)^2}}.$$

Thus,  $\operatorname{Re}(N(E))$  is the same as the describing function of the ‘average’ non-linearity  $(f_+ + f_-)/2$ .

For  $\text{Im}(N(E))$  there is a striking formula:

$$\begin{aligned}\text{Im}N(E) &= \frac{1}{\pi E} \int_{-E}^E f_+(\nu) \frac{d\nu}{E} + \frac{1}{\pi E} \int_E^{-E} f_-(\nu) \frac{d\nu}{E} \\ &= \frac{1}{\pi E^2} \int_{-E}^E [f_+(\nu) - f_-(\nu)] d\nu \\ &= -\frac{\Delta}{\pi E^2}\end{aligned}$$

where  $\Delta$  is the area enclosed by the ‘loop’ of the non-linear characteristic.

For the relay with hysteresis given above we obtain

$$\text{Re}N(E) = \frac{4R}{\pi E} \int_{\delta}^E \frac{\nu d\nu}{E^2 \sqrt{1 - (\nu/E)^2}} = \frac{4R}{\pi E} \sqrt{1 - (\delta/E)^2},$$

which gives

$$N(E) = \frac{4R}{\pi E} \left( \sqrt{1 - (\delta/E)^2} - \frac{j\delta}{E} \right).$$

### 3 Estimates of Describing Functions

Often some idea of the describing function can be obtained without calculating it in detail. The describing function is sometimes called the *equivalent linear gain*, and this name gives a pointer as to how it can be estimated. This is best illustrated by a couple of examples.

*Example: Piecewise-linear nonlinearity.* In section 2.3 we considered a nonlinearity with an incremental gain of  $m_1$  for small signals, and  $m_2$  for large signals. For a small signal the gain is exactly  $m_1$ , because the nonlinearity does not come into play. For extremely large signals the gain is essentially  $m_2$ . Clearly the effective gain for any input signal lies between these two. Thus the describing function (being real in this case) must always lie between  $m_1$  and  $m_2$ .

*Example: Relay with dead-zone.* Consider the following behaviour:

$$f(e) = \begin{cases} -1, & \text{if } e \leq -\delta \\ 0, & \text{if } |e| < \delta \\ +1, & \text{if } e \geq \delta \end{cases}$$

For  $|e| < \delta$  the describing function is clearly 0. As  $|e|$  increases slightly above  $\delta$  the output suddenly jumps to amplitude 1, so the effective gain increases quickly, and hence the describing function increases quickly (to a value which has to be calculated — see Examples Paper). As  $|e|$  increases further, the output amplitude does not increase, and so the effective gain — and hence the describing function — decreases gradually to 0.

## 4 Rigour

The first part of the describing function method, namely the prediction of limit cycles, can be made rigorous using appropriate bounds on  $g(j\omega)$  and possibly by bringing in some higher harmonics. A rather mathematical discussion along these lines and sufficient conditions for validity can be found in A.I. Mees, *Dynamics of Feedback Systems*, Wiley, 1981, Chapter 5. (This is beyond the scope of this module).

The prediction of stability/instability has so far proved difficult to treat rigorously and provide useful sufficient conditions for validity. It therefore remains as a quick method which gives a hint as to what might happen.

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Revised by J.M. Maciejowski

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